# A relaxed Beckmann's problem for SPD valued measures and application to Full Waveform Inversion

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#### Gabriele Todeschi

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**Objective**: reconstruct the velocity model  $v : D \subset \mathbb{R}^d \to \mathbb{R}_+$  of the acoustic wave equation

$$\frac{1}{v}\partial_t^2 p - \Delta p = s, \quad \text{in } [0, T] \times D$$

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Variational problem:

$$v \in \operatorname*{arginf}_{v} \mathcal{D}(d_{pred}[v], d_{obs}),$$

where D is a misfit function between true measurements  $d_{obs}$  and predicted data  $d_{pred}[v]$ .

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Measurements: pressure  $p(\cdot, x_r)$  or particles velocities  $u(\cdot, x_r) = (u_i(\cdot, x_r))_{i=1}^n$  for  $x_r \in D$ .

D model space,  $\Omega = \{x_r, .., x_{N_r}\} \times [0, T]$  data space

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The whole temporal evolution is used for the calibration  $\rightarrow$  Full Waveform Inversion Highly non-convex problem (cycle skipping)

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Optimal transport can alleviate non-convexity

Issue: OT defined for probability measures  $\longrightarrow$  need to handle the data

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 $\begin{array}{l} {\sf Kantorovich-Rubenstein}\,\,({\sf KR})\,\,{\sf norm}\\ {\sf For}\,\,\mu,\nu\in\mathcal{M}(\Omega),\,\lambda\in\mathbb{R}_+\colon\\ {\sf KR}(\mu,\nu)=\sup_{\phi}\left\{\int_{\Omega}\phi(\mu-\nu)\,,\,\,|\nabla\phi|\leq1, |\phi|\leq\lambda\right\}\end{array}$ 

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"Relaxed" transport: if  $\mu, \nu \in \mathcal{P}(\Omega)$  and  $\lambda = +\infty \implies KR(\mu, \nu) = W_1(\mu, \nu)$ 

 $\mathcal{D} = \mathit{KR}(\mu, \nu) ext{ where } \mu = \mathit{d}_{\mathsf{pred}}[\nu], \nu = \mathit{d}_{\mathsf{obs}} \longrightarrow \mathsf{good results}^1$ 

<sup>1</sup>Métivier et al., 2016

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If 
$$\int_{\Omega} \mu = \int_{\Omega} \nu$$
 and  $\lambda \gg 1$  then  
 $KR(\mu, \nu) = KR((\mu - \nu)_+, (\mu - \nu)_-) = W_1((\mu - \nu)_+, (\mu - \nu)_-)$ 

⇒ not convex with respect to shifts

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Multi-component data  $u(\cdot, x_r)$  and  $L : \mathbb{R}^n \to \mathbb{S}^n_+$ 

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Pauli's transformation (d = 2):

$$u = (u_x, u_z) \longmapsto \begin{bmatrix} lpha - u_x & u_z \\ u_z & lpha + u_x \end{bmatrix}$$
, with  $lpha = \sqrt{u_x^2 + u_z^2}$ .

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 $\mu = L(u_{\textit{pred}}[v]), \nu = L(u_{\textit{obs}}[v]) \in \mathcal{M}(\Omega; \mathbb{S}^n_+)$ 

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#### **Relaxed Beckmann's problem**

For  $\mu, \nu \in \mathcal{M}(\Omega; V_{+}), \lambda \in \mathbb{R}_{+}$ :  $\inf_{\substack{\sigma \in \mathcal{M}(\Omega; V^{n}) \\ \delta \in \mathcal{M}(\Omega; V)}} \left\{ \int_{\Omega} |\sigma|_{V^{d}} + \lambda \int_{\Omega} |\delta|_{V}, \operatorname{div}_{V}(\sigma) = \mu - \nu + \delta \right\}$ 

$$\begin{split} V &= \mathbb{S}^n, V_+ = \mathbb{S}^n_+, V^d = [\mathbb{S}^n]^d, \\ |\cdot|_V \text{ Frobenius norm, } |\cdot|_{V^d} &= \sum_{k=1}^d |\cdot|_V \\ \operatorname{div}_V : \sigma &\mapsto \sum_{i=k}^d \frac{\partial \sigma_k}{\partial x_k} \text{ and the constraint is to be considered in weak form} \end{split}$$

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#### **Relaxed Beckmann's problem**

For 
$$\mu, \nu \in \mathcal{M}(\Omega; V_{+}), \ \lambda \in \mathbb{R}_{+}$$
:  
$$\mathcal{T}^{\lambda}_{p,q}(\mu, \nu) = \inf_{\substack{\sigma \in \mathcal{M}(\Omega; V^{n}) \\ \delta \in \mathcal{M}(\Omega; V)}} \left\{ \frac{1}{p} \int_{\Omega} |\sigma|_{V^{n}}^{p} + \frac{\lambda}{q} \int_{\Omega} |\delta|_{V}^{q}, \ \mathsf{div}_{V}(\sigma) = \mu - \nu + \delta \right\}$$

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Remark: the difference of mass can be quite big due to the difference of energy

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Choice of  $p, q \in \{1, 2\}$ :

- $\circ$  Sensitivity to (small) displacements (optimal transport problem)  $\longrightarrow p = 1$

# Tuning of the model

Remark: the difference of mass can be quite big due to the difference of energy

Choice of  $p, q \in \{1, 2\}$ :

- Sensitivity to (small) displacements (optimal transport problem)  $\longrightarrow p = 1$
- Penalization proportional to the energy of the signal  $\longrightarrow q=2$

Choice of  $\lambda$ :

- For small values,  $\mathcal{T}^{\lambda}_{p,q}$  behaves like an  $L^q$  distance of  $\mu \nu$
- For big values,  $\mathcal{T}_{p,q}^{\lambda}$  behaves like an  $L^{q}$  distance on "the mass difference"
- Dimensional analysis:

$$[\mu] = [\nu] = A, \ [\sigma] = A \cdot I \quad \longrightarrow \quad [\lambda] = A^{p-q} \cdot I^p$$

where I is a length, A an amplitude

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Dual structure is useful for exact solutions, derivatives, numerical computation,...

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$$egin{aligned} p &= 1, q = 1 \ && \mathcal{T}_{1,1}^\lambda(\mu,
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 $\longrightarrow\,$  vectorial extension of the KR norm

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$$egin{aligned} m{p} &= 1, m{q} = 2 \ \mathcal{T}_{1,2}^\lambda(\mu,
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u-\mu >_V -rac{1}{2\lambda} \int_\Omega |\phi|_V^2 \,, \; |
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 $\longrightarrow$  quadratic penalization of the potential in the KR

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 $\longrightarrow\,$  quadratic penalization of the potential in the KR

$$\begin{split} \rho &= 2, q = 2 \\ \mathcal{T}_{2,2}^{\lambda}(\mu,\nu) &= \sup_{\phi} -\frac{1}{2} \int_{\Omega} |\nabla \phi|_{V^n}^2 - \frac{1}{2\lambda} \int_{\Omega} |\phi|_{V}^2 + \int_{\Omega} \langle \phi, \nu - \mu \rangle_{V} \end{split}$$

 $\rightarrow$  not a transport, not coupled...

Consider two delta measures  $\mu = M_1 \delta_{x_1}, \nu = M_2 \delta_{x_2}$ ,  $M_1, M_2 \in \mathbb{S}^n_+$ 

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 $\longrightarrow$  (generalized) Fermat-Torricelli problem

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 $\longrightarrow\,\lambda$  represents (half) the distance at which mass is transported

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The optimal potential is:  $\phi(x_1) = a, \phi(x_2) = b$ 

$$\begin{cases} a = \sqrt{\lambda(m_1 - m_2)}, b = a - |x_1 - x_2| & \text{if } \lambda \ge \frac{|x_1 - x_2|^2}{m_1 - m_2} \\ a = \frac{|x_1 - x_2|}{2} + \frac{\lambda(m_1 - m_2)}{2|x_1 - x_2|}, b = -\frac{|x_1 - x_2|}{2} + \frac{\lambda(m_1 - m_2)}{2|x_1 - x_2|} & \text{if } \lambda \le \frac{|x_1 - x_2|^2}{|m_1 - m_2|} \\ a = \sqrt{\lambda m_1}, b = -\sqrt{\lambda m_2} & \text{if } \lambda \le \frac{|x_1 - x_2|^2}{(\sqrt{m_1} + \sqrt{m_2})^2} \end{cases}$$

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 $\longrightarrow$   $\lambda$  represents the distance at which a unit of mass is transported

The tuning of  $\lambda$  can be done by choosing a reference transport distance and a reference mass to be transported

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## Numerical solution

Numerical solution of  $\mathcal{T}_{p,q}^{\lambda}$ :

- Finite difference discretization
- SDMM algorithm: primal-dual proximal splitting technique
- Bottleneck: solution of a Poisson equation
- The SPD transport does not reflect on the complexity of the algorithm
- $(p = 1, q = 2) \rightarrow$  the higher regularity can be exploited (FISTA, Chambolle-Pock, ...)

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Solution of the inverse problem:

- I-BFGS algorithm: requires  $\frac{\partial \mathcal{D}(\mu[v],\nu)}{\partial v}$
- Derivative via adjoint state method: requires  $rac{\partial \mathcal{T}^{\lambda}_{
  ho,q}(\mu,
  u)}{\partial \mu}$
- thanks to the dual structure:

$$\frac{\partial \mathcal{T}_{p,q}^{\lambda}(\mu,\nu)}{\partial \mu} = \phi$$

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#### Numerical results: sensitivity to time shift



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 $D = [a, b] \times [c, d] \subset \mathbb{R}^2, n = 2, T = 3000, N_r = 169$ 

 $\Omega = \{x_1, .., x_{N_r}\} \times [0, T] \text{ semi-discrete 2d domain}$ 

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"Dimensional anisotropy":  $\tilde{\Omega}=\{x_r,..,x_{N_r}\}\times[0,\,T/\bar{v}]$  where  $\bar{v}$  mean velocity

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# Conclusions

We introduced an (unbalanced) L1 transport for SPD valued measures to transport multi-component signals

Computationally affordable

Good results for the FWI problem justifying the approach

Applications to other problems can be foreseen

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# Conclusions

We introduced an (unbalanced) L1 transport for SPD valued measures to transport multi-component signals

Computationally affordable

Good results for the FWI problem justifying the approach

Applications to other problems can be foreseen

BUT:

- Rely on a lift function: linked to the physics of the problem at hand
- Sensitivity to the calibration of  $\lambda$  (depending again on the lift/physics)

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# Thank you for your attention!

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